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# COMPARISON OF EIGENPROBLEM SOLUTION APPROACHES FOR DIRICHLET CONDITIONS IN IRREGULAR DOMAINS VIA INTEGRAL TRANSFORMS

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**Abstract.** This paper is focused on the solution of eigenvalue problems for irregular-shaped domains using the Generalized Integral Transform Technique (GITT). Two-dimensional Helmholtz problems in cartesian coordinates with Dirichlet boundary conditions are used as test-cases with a semi-circular region with angles  $\phi = 90^\circ$  and  $\phi = 180^\circ$ . The solution of the eigenproblem is carried out by using an auxiliary problem constructed from simpler 1D eigenfunctions as a basis for expanding the sought solution, and two-different approaches are comparative analyzed. The first is based on using an auxiliary problem defined within the same region as the original problem (the coincident domain approach – CDA), while the second relies on using an auxiliary problem defined in a fictitious domain, obtained from expanding the original domain to a larger regular domain (the fictitious domain approach – FDA). The solutions are verified by comparing the calculated results of the semi-circular domain with the exact analytical solution obtained by employing cylindrical coordinates.

**Keywords:** Irregular Domains, GITT, Eigenproblem, Symbolical Computation

## NOMENCLATURE

|                      |                          |                    |                                   |
|----------------------|--------------------------|--------------------|-----------------------------------|
| $A, B, C, M$         | integral coefficients    | $\mu$              | eigenvalues                       |
| $k$                  | coefficient              | $\Psi$             | eigenfunctions                    |
| $w$                  | weighting function       | $\Omega$           | basis of auxiliary eigenfunctions |
| $d, G$               | known functions          | $\phi$             | angle                             |
| $\mathcal{A}$        | area                     | <b>Subscripts</b>  |                                   |
| $\mathcal{S}$        | surface                  | $i, j, k, l, m, n$ | summation indices                 |
| $\mathcal{V}$        | volume                   | c                  | circle                            |
| $x, y$               | spatial coordinates      | fic                | fictitious                        |
| $X, Y$               | auxiliary eigenfunctions | max                | maximum                           |
| <b>Greek Symbols</b> |                          | <b>Overscripts</b> |                                   |
| $\beta, \omega$      | auxiliary eigenvalues    | –                  | transformed function              |
| $\delta$             | Kronecker delta function |                    |                                   |

## 1. INTRODUCTION

Previously to the interpretation of a studied problem, many considerations need to be taken into account primarily and the choice of the working domain is a concerning subject to consider. There is no denying that an abundant number of research papers in the literature deals with different problems in traditional regular geometries, however most engineering applications are actually formulated for irregular geometries. Based on that, when irregular domains are treated, a numerical methodology is required to obtain the solution of eigenproblems. Due to the hybrid analytical-numerical nature of the Generalized Integral Transform Technique (GITT) (Cotta, 1993), this method seems a convenient choice to the

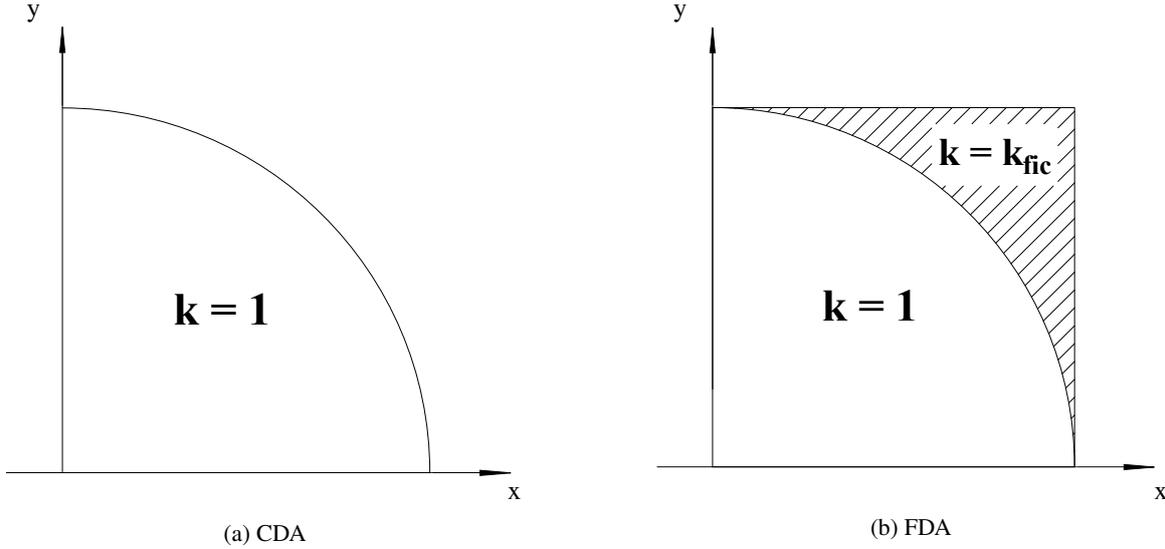


Figure 1: Representation of both approaches in semi-circular domain.

development of the eigenvalue problem applied to irregular-shaped domains.

The use of the GITT method is a fairly common methodology applied to regular geometries, as different applications of the method can be cited, for example, studies presenting the GITT for thermally developing flows in rectangular ducts (Chalhub and Sphaier, 2010; Aparecido and Cotta, 1990), for forced convection analysis in rectangular cross-section (Lindquist and Aparecido, 1999) and slip-flow in microchannels (Yu and Ameen, 2001). However, even sparse, some investigations into irregular domains can also be presented, as in fully developed flows (Barbuto and Cotta, 1997; Pérez Guerrero *et al.*, 2000), fluid flow in irregularly-shaped microchannels (Pinheiro and Sphaier, 2016) or in specific studies into eigenproblems in irregular-shaped geometries (Sphaier and Cotta, 2000).

The steps for the GITT approach requires the solution of the Sturm-Liouville eigenvalue problem, which determines the set of eigenfunctions, related to corresponding eigenvalues. In the literature, the eigenproblems can incorporate primarily importance in specific works, such as the one presented by Cotta *et al.* (2016), in which nonlinear convection-diffusion problems yield also nonlinear eigenvalue problem solved by means of the GITT method. Early works involving eigenproblems can be exemplified by the work proposed by Tai and Shaw (1974), where the two or three-dimensional Helmholtz-based eigenvalue problem is numerically solved in arbitrarily-shaped domains.

As many engineering problems are indeed applied for irregular or more complex geometries, this study presents two different strategies to solve eigenproblems in irregularly-shaped domains, in which Dirichlet boundary conditions are considered. One approach is based on the accordance of the actual geometry and the domain related to the mathematical formulation. This approach is named the Coincident Domain Approach (CDA). Regarding the second approach, the Fictitious Domain Approach (FDA) requires the definition of a regular domain outside the irregular geometry, characterized by a fictitious coefficient  $k_{fic}$ . Both approaches (figure 1) are solved via GITT with the aid of the Mathematica Platform (Wolfram, 2003) and a convergence analysis is employed with a semi-circular region.

## 2. MATHEMATICAL FORMULATION

The Eigenvalue problem is represented by the Sturm-Liouville Equation, in which the eigenvalues are illustrated by  $\mu$  and the corresponding eigenfunctions are characterized by  $\Psi$ . The generalized eigenproblem is established for an arbitrary domain  $\mathcal{V}$  with delimited surface  $\mathcal{S}$ , as can be seen in the following equations:

$$\nabla \cdot (k(\mathbf{x}) \nabla \Psi) + (\mu^2 w(\mathbf{x}) - d(\mathbf{x})) \Psi = 0, \quad \mathbf{x} \in \mathcal{V}, \quad (1)$$

$$k(\mathbf{x}) (\nabla \Psi \cdot \mathbf{n}) + \gamma(\mathbf{x}) \Psi = 0, \quad \mathbf{x} \in \mathcal{S}, \quad (2)$$

where  $w(\mathbf{x})$  is considered a weighting function and  $d(\mathbf{x})$  is known function. In order to apply the GITT approach, the auxiliary transform-inverse pair is described:

$$\Psi_n = \sum_{i=1}^{\infty} \Omega_i(\mathbf{x}) \bar{\Psi}_i^n, \quad (3)$$

$$\bar{\Psi}_i^n = \int_{\mathcal{V}^*} w^*(\mathbf{x}) \Psi_n(\mathbf{x}) \Omega_i(\mathbf{x}) d\mathcal{V} \quad (4)$$

The auxiliary eigenfunctions are orthogonal, so:

$$\int_{\mathcal{V}^*} w^*(\mathbf{x}) \Omega_i(\mathbf{x}) \Omega_j(\mathbf{x}) d\mathcal{V} = \delta_{i,j}, \quad (5)$$

in which  $\delta_{i,j}$  is the Kronecker delta function and  $w^*(\mathbf{x})$  is the weighting function, now associated with functions  $\Omega$ .

The procedure to transform the original eigenproblem requires the multiplication of the differential equation (1) by the auxiliary eigenfunctions  $\Omega$ , integrating within the considered domain  $\mathcal{V}_T$ , delimited by surface  $\mathcal{S}_T$ . The domain can be related to  $\mathcal{V}$  or  $\mathcal{V}^*$ , depending only on the strategy applied (Coincident Domain Approach or Fictitious Domain Approach).

$$\int_{\mathcal{V}_T} \nabla \cdot (k \nabla \Psi) \Omega_i d\mathcal{V} + \mu^2 \int_{\mathcal{V}_T} w \Psi \Omega_i d\mathcal{V} - \int_{\mathcal{V}_T} d \Psi \Omega_i d\mathcal{V} = 0 \quad (6)$$

In the first integral, the second Green formula is used, which gives:

$$\int_{\mathcal{V}_T} \nabla \cdot (k \nabla \Omega_i) \Psi d\mathcal{V} + \int_{\mathcal{S}_T} k (\Omega_i \nabla \Psi - \Psi \nabla \Omega_i) \cdot \mathbf{n} d\mathcal{A} + \mu^2 \int_{\mathcal{V}_T} w \Psi \Omega_i d\mathcal{V} - \int_{\mathcal{V}_T} d \Psi \Omega_i d\mathcal{V} = 0 \quad (7)$$

And the first Green formula gives:

$$- \int_{\mathcal{V}_T} (\nabla \Omega_i) \cdot (k \nabla \Psi) d\mathcal{V} + \int_{\mathcal{S}_T} [k (\Omega_i \nabla \Psi)] \cdot \mathbf{n} d\mathcal{A} + \mu^2 \int_{\mathcal{V}_T} w \Psi \Omega_i d\mathcal{V} - \int_{\mathcal{V}_T} d \Psi \Omega_i d\mathcal{V} = 0 \quad (8)$$

Both first and second Green identities can be applied for the considered problem, however when discontinuous regions are in place, the first Green formula is recommended. The same apply for the FDA, because the first Green formula results in simpler coefficients to be calculated.

Applying the Dirichlet boundary condition on the irregular domain surface, one can infer that:

$$\Psi = 0 \quad \text{and} \quad \Omega_i = 0 \quad \text{in} \quad \mathbf{x} \in \mathcal{S}_T \quad (9)$$

which reduces the surface integrals of equations (7) and (8) to zero.

After substituting the integrals for coefficients  $A$ ,  $B$ ,  $C$  and  $D$ , the final transformed eigenproblem for both first and second Green identities are given by:

$$\sum_{j=1}^{\infty} (-D_{i,j} + \mu^2 B_{i,j} - C_{i,j}) \bar{\Psi}_j = 0 \quad (10)$$

$$\sum_{j=1}^{\infty} (A_{i,j} + \mu^2 B_{i,j} - C_{i,j}) \bar{\Psi}_j = 0 \quad (11)$$

where  $A$ ,  $B$ ,  $C$  and  $D$  are specifically:

$$A_{i,j} = \int_{\mathcal{V}_T} \nabla \cdot (k \nabla \Omega_j) \Omega_i d\mathcal{V} \quad (12)$$

$$B_{i,j} = \int_{\mathcal{V}_T} w \Omega_j \Omega_i d\mathcal{V} \quad (13)$$

$$C_{i,j} = \int_{\mathcal{V}_T} d \Omega_j \Omega_i d\mathcal{V} \quad (14)$$

$$D_{i,j} = \int_{\mathcal{V}_T} (\nabla \Omega_i) \cdot (k \nabla \Omega_j) d\mathcal{V}, \quad (15)$$

Another way to write the proposed solution can be seen below:

$$\sum_{j=1}^{\infty} (M_{i,j} + \mu^2 B_{i,j}) \bar{\Psi}_j = 0 \quad (16)$$

where the integral coefficient  $M$  is the summation of both  $A$  and  $C$  or  $D$  and  $C$ , as observed:

$$M_{i,j} = A_{i,j} - C_{i,j} \quad (17)$$

$$M_{i,j} = -D_{i,j} - C_{i,j} \quad (18)$$

The system is then truncation to a finite number of terms, giving the eigenvalues ( $\mu$ ) and the eigenvectors ( $\bar{\Psi}$ ).

Regarding the considering domain, when discussing the CDA, it requires that the domain and the actual considered geometry be coincident, in which  $\mathcal{S} = \mathcal{S}^* = \mathcal{S}_T$  and  $\mathcal{V} = \mathcal{V}^* = \mathcal{V}_T$ . However, in the FDA, although  $\mathcal{V}^* = \mathcal{V}_T$ , the problem is extended to fictitious domain, which included the irregular domain, in other words  $\mathcal{V} \subset \mathcal{V}_T$ . For the considered test-case, the extended domain is considered a rectangular geometry, which includes the semi-circular region.

### 3. TEST-CASE

For the considered test-cases with both strategies, the differential equation is based on a two-dimensional Helmholtz eigenproblem, as can be seen:

$$\frac{\partial^2 \Psi}{\partial x^2} + \frac{\partial^2 \Psi}{\partial y^2} + \mu^2 \Psi = 0, \quad 0 \leq y \leq y_1(x), \quad 0 \leq x \leq 1, \quad (19)$$

$$\Psi(x, 0) = 0, \quad 0 \leq x \leq 1, \quad (20)$$

$$\Psi(0, y) = 0, \quad 0 \leq y \leq y(0), \quad (21)$$

$$\Psi(0, y(x)) = 0, \quad y = y(x), \quad 0 \leq x \leq 1, \quad (22)$$

where auxiliary eigenfunctions are determined by:

$$\Omega_i(x, y) = X_n(x) Y_k(y), \quad (23)$$

in which the solutions for  $X_n$  and  $Y_k$  are obtained by the following one-dimensional Helmholtz problems:

$$\frac{\partial^2 Y_k}{\partial y^2} + \omega_k^2 Y_k = 0, \quad (24)$$

$$\frac{d^2 X_n}{dx^2} + \beta_n^2 X_n = 0, \quad (25)$$

For both CDA and FDA schemes, the boundary conditions for the 1D eigenproblems are:

$$X(0) = 0, \quad X(1) = 0, \quad Y(y = 0) = 0, \quad (26)$$

#### 3.1 CDA

The CDA is commonly seen in many problems presented in the literature. It requires that the considered region coincides with the domain applied in the mathematical formulation. For that, the boundary condition in the irregular region is represented by:

$$Y = 0, \quad \text{at } y = y_c(x) \quad (27)$$

The function  $y_c$  is the equation for the semi-circular region, which is:

$$y_c(x) = \sqrt{1 - x^2} \quad (28)$$

Considering the known function  $d$  equals zero and integral coefficients in two-dimensional cartesian coordinates, A and B are characterized by:

$$A_{i,j} = \int_0^1 \int_0^{y_c(x)} \left[ X_n \frac{\partial^2 Y_k}{\partial x^2} + 2 \frac{\partial X_n}{\partial x} \frac{\partial Y_k}{\partial x} - (\beta_n^2 + \omega_k^2) X_n Y_k \right] X_m Y_l dy dx \quad (29)$$

$$B_{i,j} = \int_0^1 \int_0^{y_c(x)} X_m Y_l X_n Y_k dy dx \quad (30)$$

The integral coefficient B is the same for both CDA and FDA. As for the integral coefficient A, it is acquired by the routine **NIntegrate** from the *Mathematica* platform.

#### 3.2 FDA

The proposed FDA scheme requires the definition of an extended domain which contains the irregular-shaped region but doesn't match it. Because of that, the Dirichlet boundary condition is to be considered:

$$Y = 0, \quad \text{at } y = 1 \quad (31)$$

which implies that the extended surface is a rectangular one. This extended domain is characterized by a coefficient  $k_{fic}$ , so the original coefficient  $k$  can be expressed by:

$$k = 1, \quad \text{if } y \leq y_c(x) \quad \text{or} \quad k = k_{fic}, \quad \text{if } y > y_c(x) \quad (32)$$

The integral coefficient D is then expressed by:

$$D_{i,j} = \int_0^1 \int_0^1 k \left[ X_n X_m \frac{\partial Y_l}{\partial y} \frac{\partial Y_k}{\partial y} + \left( Y_l \frac{\partial X_n}{\partial x} + X_n \frac{\partial Y_l}{\partial x} \right) \left( Y_k \frac{\partial X_m}{\partial x} + X_m \frac{\partial Y_k}{\partial x} \right) \right] dy dx, \quad (33)$$

which can be rewritten to avoid too many numerical integrations:

$$D_{i,j} = \int_0^1 \int_0^1 G_{k,l,n,m}(x, y) dy dx + \int_0^1 \int_{y_c(x)}^1 G_{k,l,n,m}(x, y) dy dx \quad (34)$$

$$G_{k,l,n,m} = \left[ X_n X_m \frac{\partial Y_l}{\partial y} \frac{\partial Y_k}{\partial y} + \left( Y_l \frac{\partial X_n}{\partial x} + X_n \frac{\partial Y_l}{\partial x} \right) \left( Y_k \frac{\partial X_m}{\partial x} + X_m \frac{\partial Y_k}{\partial x} \right) \right] \quad (35)$$

The calculation of this coefficient can be partially solved analytically, so the function **Integrate** from the *Mathematica* software is required. Although not all calculation can be done analytically, it can be concluded that the implementation of this coefficient has a lower computation cost when considered to the CDA scheme, which is purely numerically obtained.

#### 4. RESULTS AND DISCUSSION

One important remark before the presentation of the results is that, for both CDA and FDA, the solution will be presented in semi-circular regions with  $\phi = 90^\circ$  and  $\phi = 180^\circ$ , as described in figure 2.

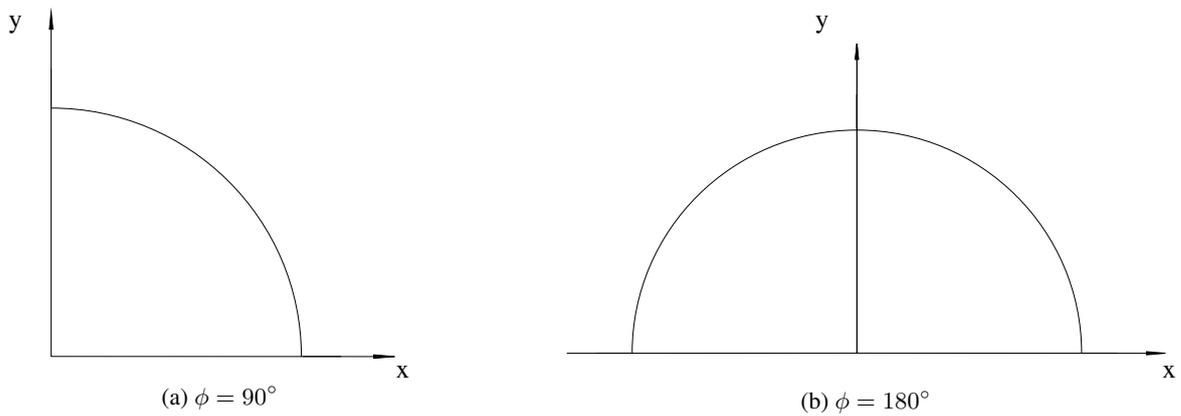


Figure 2: Considered geometries for convergence analysis.

The first set of results is presented for the semi-circular domain with the Coincident Domain Approach, as shown in table 1, where the analytical value in cylindrical coordinates is considered. Considering that the converged eigenvalue is the analytical solution in cylindrical coordinates, one can infer that a maximum of 6 fully converged digits were obtained with as low as 100 terms for some eigenvalues, showing satisfactory error estimates.

For the FDA scheme, the solution can be compared with different values of  $k_{fic}$ . In order to acquire the convergence behavior, values of  $k_{fic} = 10$  and  $100$  are investigated in tables 2 and 3. As can be seen, the convergence for the FDA is clearly worse than the CDA with a maximum of 2 fully converged digits obtained in some cases. Specifically, as higher values of  $k_{fic}$  are considered, the convergence analysis gets worse with the same number of truncation orders than for smaller values of  $k_{fic}$ . However, a clear path can be delimited: smaller values of  $k_{fic}$  converge to distant values from the analytical solution, as for higher values of  $k_{fic}$  need more terms to converge, but the values seem to get closer to the analytical ones. If a triangular would be consider, the whole implementation would be purely analytical, so more terms can be described. Although a worse convergence is obtained, the FDA has a advantage towards the CDA, which is that a irregular domain with a cavity or hole can be described only in FDA scheme.

Table 1: Convergence behavior in a semi-circular geometry with Dirichlet Boundary Conditions – CDA.

| $\phi = 90^\circ$  |         |         |         |         |          |          |          |          |          |            |
|--------------------|---------|---------|---------|---------|----------|----------|----------|----------|----------|------------|
| $k_{max}$          | $\mu_1$ | $\mu_2$ | $\mu_3$ | $\mu_4$ | $\mu_5$  | $\mu_6$  | $\mu_7$  | $\mu_8$  | $\mu_9$  | $\mu_{10}$ |
| 20                 | 5.13832 | 7.60342 | 8.42181 | 9.98437 | 11.09534 | 11.62904 | 12.35295 | 13.66862 | 14.49606 | 14.79870   |
| 40                 | 5.13652 | 7.59344 | 8.41873 | 9.95080 | 11.07236 | 11.62193 | 12.26449 | 13.61275 | 14.38323 | 14.56280   |
| 60                 | 5.13612 | 7.59124 | 8.41806 | 9.94424 | 11.06902 | 11.62097 | 12.24140 | 13.60093 | 14.37814 | 14.51187   |
| 80                 | 5.13596 | 7.59026 | 8.41780 | 9.94149 | 11.06752 | 11.62061 | 12.23631 | 13.59679 | 14.37604 | 14.49716   |
| 100                | 5.13585 | 7.58966 | 8.41762 | 9.93983 | 11.06664 | 11.62037 | 12.23280 | 13.59444 | 14.37495 | 14.48972   |
| 250                | 5.13568 | 7.58868 | 8.41734 | 9.93708 | 11.06519 | 11.61997 | 12.22716 | 13.59061 | 14.37315 | 14.47916   |
| 500                | 5.13564 | 7.58845 | 8.41728 | 9.93644 | 11.06487 | 11.61988 | 12.22580 | 13.58974 | 14.37275 | 14.47680   |
| 750                | 5.13563 | 7.58840 | 8.41726 | 9.93628 | 11.06479 | 11.61986 | 12.22546 | 13.58952 | 14.37264 | 14.47618   |
| 1000               | 5.13563 | 7.58838 | 8.41725 | 9.93622 | 11.06476 | 11.61986 | 12.22533 | 13.58944 | 14.37261 | 14.47594   |
| $\infty$           | 5.13562 | 7.58834 | 8.41724 | 9.93611 | 11.06471 | 11.61984 | 12.22509 | 13.58892 | 14.37254 | 14.47550   |
| $\phi = 180^\circ$ |         |         |         |         |          |          |          |          |          |            |
| $k_{max}$          | $\mu_1$ | $\mu_2$ | $\mu_3$ | $\mu_4$ | $\mu_5$  | $\mu_6$  | $\mu_7$  | $\mu_8$  | $\mu_9$  | $\mu_{10}$ |
| 20                 | 3.83329 | 5.14647 | 6.40266 | 7.01908 | 7.64424  | 8.44744  | 8.85079  | 9.79459  | 10.21495 | 10.24946   |
| 40                 | 3.83238 | 5.13849 | 6.38917 | 7.01688 | 7.60372  | 8.42201  | 8.80449  | 9.77749  | 9.98520  | 10.17626   |
| 60                 | 3.83203 | 5.13719 | 6.38482 | 7.01619 | 7.59711  | 8.41983  | 8.78782  | 9.76861  | 9.96227  | 10.17441   |
| 80                 | 3.83192 | 5.13672 | 6.38293 | 7.01598 | 7.59445  | 8.41906  | 8.78148  | 9.76539  | 9.95280  | 10.17406   |
| 100                | 3.83187 | 5.13632 | 6.38231 | 7.01588 | 7.59229  | 8.41840  | 8.77911  | 9.76435  | 9.94730  | 10.17391   |
| 250                | 3.83174 | 5.13580 | 6.38066 | 7.01565 | 7.58935  | 8.41753  | 8.77332  | 9.76178  | 9.93900  | 10.17356   |
| 500                | 3.83172 | 5.13568 | 6.38033 | 7.01561 | 7.58869  | 8.41734  | 8.77211  | 9.76127  | 9.93711  | 10.17350   |
| 750                | 3.83171 | 5.13565 | 6.38025 | 7.01560 | 7.58852  | 8.41729  | 8.77181  | 9.76116  | 9.93663  | 10.17348   |
| 1000               | 3.83171 | 5.13564 | 6.38022 | 7.01559 | 7.58846  | 8.41728  | 8.77169  | 9.76111  | 9.93645  | 10.17348   |
| $\infty$           | 3.83171 | 5.13562 | 6.38016 | 7.01559 | 7.58834  | 8.41724  | 8.77148  | 9.76102  | 9.93611  | 10.17350   |

Table 2: Convergence in a semi-circular geometry with  $\phi = 90^\circ$  – FDA.

| $\phi = 90^\circ$ |         |         |         |          |          |          |          |          |          |            |
|-------------------|---------|---------|---------|----------|----------|----------|----------|----------|----------|------------|
| $k = 10$          |         |         |         |          |          |          |          |          |          |            |
| $k_{max}$         | $\mu_1$ | $\mu_2$ | $\mu_3$ | $\mu_4$  | $\mu_5$  | $\mu_6$  | $\mu_7$  | $\mu_8$  | $\mu_9$  | $\mu_{10}$ |
| 20                | 5.27685 | 7.83851 | 8.67109 | 10.29226 | 11.48134 | 12.04082 | 12.78651 | 14.19035 | 15.10643 | 15.52958   |
| 40                | 5.21015 | 7.74714 | 8.53850 | 10.14279 | 11.30977 | 11.78697 | 12.52116 | 13.87470 | 14.72583 | 14.87744   |
| 60                | 5.17914 | 7.70698 | 8.48266 | 10.08553 | 11.24295 | 11.70077 | 12.44817 | 13.78430 | 14.61807 | 14.75799   |
| 80                | 5.16458 | 7.68162 | 8.45777 | 10.04950 | 11.20208 | 11.66363 | 12.39344 | 13.72910 | 14.55695 | 14.70063   |
| 100               | 5.15174 | 7.66188 | 8.43584 | 10.02730 | 11.17103 | 11.63152 | 12.36431 | 13.69554 | 14.51301 | 14.66209   |
| 250               | 5.12034 | 7.61689 | 8.38361 | 9.96468  | 11.10302 | 11.55888 | 12.27967 | 13.60807 | 14.42004 | 14.55194   |
| 500               | 5.10395 | 7.59215 | 8.35679 | 9.93259  | 11.06636 | 11.52218 | 12.23905 | 13.56397 | 14.37164 | 14.50103   |
| 750               | 5.09650 | 7.58158 | 8.34466 | 9.91922  | 11.05079 | 11.50580 | 12.22161 | 13.54584 | 14.35137 | 14.48092   |
| 1000              | 5.09200 | 7.57489 | 8.33734 | 9.91077  | 11.04099 | 11.49590 | 12.21107 | 13.53439 | 14.33865 | 14.46838   |
| $\infty$          | 5.13562 | 7.58834 | 8.41724 | 9.93611  | 11.06471 | 11.61984 | 12.22509 | 13.58892 | 14.37254 | 14.47550   |
| $k = 100$         |         |         |         |          |          |          |          |          |          |            |
| $k_{max}$         | $\mu_1$ | $\mu_2$ | $\mu_3$ | $\mu_4$  | $\mu_5$  | $\mu_6$  | $\mu_7$  | $\mu_8$  | $\mu_9$  | $\mu_{10}$ |
| 20                | 5.58126 | 8.29483 | 9.20839 | 10.89378 | 12.25162 | 12.89861 | 13.54652 | 15.46478 | 19.50706 | 22.92077   |
| 40                | 5.44814 | 8.05164 | 8.94860 | 10.55835 | 11.77367 | 12.39922 | 13.02801 | 14.50640 | 15.38266 | 15.50791   |
| 60                | 5.40333 | 7.96057 | 8.86742 | 10.43190 | 11.62541 | 12.26992 | 12.84309 | 14.29839 | 15.13977 | 15.25360   |
| 80                | 5.35507 | 7.91061 | 8.78340 | 10.34984 | 11.54597 | 12.13932 | 12.75049 | 14.17143 | 15.02073 | 15.11774   |
| 100               | 5.33305 | 7.87407 | 8.74540 | 10.30549 | 11.48857 | 12.08280 | 12.68417 | 14.10413 | 14.93930 | 15.05078   |
| 250               | 5.26019 | 7.77032 | 8.62257 | 10.16860 | 11.33202 | 11.90571 | 12.50861 | 13.90984 | 14.72359 | 14.81602   |
| 500               | 5.22092 | 7.71559 | 8.55740 | 10.09871 | 11.25092 | 11.81403 | 12.42590 | 13.81256 | 14.61565 | 14.71291   |
| 750               | 5.20426 | 7.69210 | 8.52990 | 10.06859 | 11.21637 | 11.77568 | 12.38903 | 13.77096 | 14.57019 | 14.66922   |
| 1000              | 5.19389 | 7.67743 | 8.51282 | 10.05025 | 11.19484 | 11.75194 | 12.36615 | 13.74568 | 14.54198 | 14.64291   |
| $\infty$          | 5.13562 | 7.58834 | 8.41724 | 9.93611  | 11.06471 | 11.61984 | 12.22509 | 13.58892 | 14.37254 | 14.47550   |

Table 3: Convergence in a semi-circular geometry with  $\phi = 180^\circ$  – FDA.

| $\phi = 180^\circ$ |         |         |         |         |         |          |          |          |          |            |
|--------------------|---------|---------|---------|---------|---------|----------|----------|----------|----------|------------|
| $k = 10$           |         |         |         |         |         |          |          |          |          |            |
| $k_{max}$          | $\mu_1$ | $\mu_2$ | $\mu_3$ | $\mu_4$ | $\mu_5$ | $\mu_6$  | $\mu_7$  | $\mu_8$  | $\mu_9$  | $\mu_{10}$ |
| 20                 | 4.00924 | 5.34352 | 6.75114 | 7.40540 | 8.06183 | 8.84029  | 9.40971  | 10.78878 | 10.83126 | 12.34792   |
| 40                 | 3.92992 | 5.28781 | 6.56314 | 7.20502 | 7.88639 | 8.69281  | 9.04732  | 10.07393 | 10.38591 | 10.47497   |
| 60                 | 3.90923 | 5.23346 | 6.51335 | 7.16179 | 7.79084 | 8.58228  | 8.98431  | 9.97501  | 10.21098 | 10.39627   |
| 80                 | 3.89664 | 5.21272 | 6.48971 | 7.13668 | 7.75720 | 8.54302  | 8.93693  | 9.93103  | 10.14743 | 10.35489   |
| 100                | 3.88688 | 5.19470 | 6.46981 | 7.11770 | 7.72934 | 8.51068  | 8.90750  | 9.89606  | 10.12390 | 10.32479   |
| 250                | 3.85515 | 5.14395 | 6.42022 | 7.05725 | 7.65206 | 8.42271  | 8.83362  | 9.81425  | 10.01184 | 10.23292   |
| 500                | 3.83893 | 5.12042 | 6.39220 | 7.02730 | 7.61729 | 8.38372  | 8.79482  | 9.76945  | 9.96568  | 10.18960   |
| 750                | 3.83148 | 5.11059 | 6.37971 | 7.01364 | 7.60193 | 8.36763  | 8.77724  | 9.74989  | 9.94612  | 10.17002   |
| 1000               | 3.82662 | 5.10438 | 6.37175 | 7.00475 | 7.59323 | 8.35748  | 8.76648  | 9.73740  | 9.93393  | 10.15743   |
| $\infty$           | 3.83171 | 5.13562 | 6.38016 | 7.01559 | 7.58834 | 8.41724  | 8.77148  | 9.76102  | 9.93611  | 10.17350   |
| $k = 100$          |         |         |         |         |         |          |          |          |          |            |
| $k_{max}$          | $\mu_1$ | $\mu_2$ | $\mu_3$ | $\mu_4$ | $\mu_5$ | $\mu_6$  | $\mu_7$  | $\mu_8$  | $\mu_9$  | $\mu_{10}$ |
| 20                 | 4.31085 | 5.84473 | 7.25384 | 8.04603 | 8.59306 | 10.13260 | 11.94728 | 14.79482 | 17.68540 | 19.92979   |
| 40                 | 4.16520 | 5.58693 | 6.97770 | 7.66920 | 8.29973 | 9.22093  | 9.66237  | 10.78490 | 11.05393 | 11.29099   |
| 60                 | 4.10267 | 5.53623 | 6.83554 | 7.53353 | 8.15258 | 9.11942  | 9.39369  | 10.50216 | 10.70373 | 10.99601   |
| 80                 | 4.06463 | 5.47314 | 6.77626 | 7.45508 | 8.10143 | 8.99552  | 9.33970  | 10.39134 | 10.58283 | 10.85016   |
| 100                | 4.03529 | 5.41877 | 6.72310 | 7.39714 | 7.99636 | 8.89577  | 9.25475  | 10.30210 | 10.47652 | 10.75046   |
| 250                | 3.95842 | 5.31570 | 6.58937 | 7.24952 | 7.85385 | 8.71591  | 9.06233  | 10.08458 | 10.27173 | 10.51795   |
| 500                | 3.92011 | 5.26139 | 6.52958 | 7.17814 | 7.77259 | 8.62455  | 8.97654  | 9.99075  | 10.17218 | 10.41100   |
| 750                | 3.90323 | 5.23614 | 6.49950 | 7.14689 | 7.73726 | 8.58258  | 8.93721  | 9.94420  | 10.12641 | 10.36471   |
| 1000               | 3.89342 | 5.22291 | 6.48320 | 7.12880 | 7.71765 | 8.56069  | 8.91451  | 9.91902  | 10.10209 | 10.33816   |
| $\infty$           | 3.83171 | 5.13562 | 6.38016 | 7.01559 | 7.58834 | 8.41724  | 8.77148  | 9.76102  | 9.93611  | 10.17350   |

## 5. CONCLUSIONS

This paper proposed two schemes to acquire eigenproblem solutions in irregularly-shaped domains, using a semi-circular region as test-case. As a general representation of the problem is defined for both strategies, specific formulations were developed for the considered test-case. The CDA and the FDA were implemented based on the Dirichlet Boundary Conditions and a convergence investigation was conducted for both approaches.

As the results were presented, one important trend observed was that the CDA scheme presented itself as a better fit for the considered test-case within a semi-circular region in cartesian coordinates. However, the proposed methodology for CDA cannot be exploited for region with orifices, as the FDA methodology can. One strategy to consider to improve the solution methodology for the FDA requires that the integral coefficient could be calculated via analytical implementation, which would result in more terms to be consider and better error estimates for the FDA.

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## 8. RESPONSIBILITY NOTICE

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